# **Turbulence Induced Dust Collision Velocities and**





## Rates

Alexander Hubbard





#### **PROBLEM:**

Early dust growth in protoplanetary disks occurs collisionally, and, for a broad range of dust sizes, is driven by turbulent flows. The rate of growth, and the final dust grain size distribution, depend sensitively on the velocities and rates of these collisions.

### **EQUATIONS:**

**Particles** are entrained by gas drag with stopping time  $\mathbf{T}_{\mathbf{p}}$ :  $\Box$ 

 $\partial \boldsymbol{u}_p(t)$  $\boldsymbol{u}_p(t) - \boldsymbol{V}(\boldsymbol{x}_p, t)$  $\tau_p \leftarrow \text{Stopping}$  time

#### **POINT PARTICLES:**

As a further complication, dust grains in protoplanetary disks are tiny compared with the turbulent motions. They must be treated as point particles. To be able to make statements about the turbulence induced collisions we need to be able to study the turbulence induced collision parameters in the limit of dust grain separation going to zero.

**Estimating** the velocity scale at which turbulence induced collisions occur is straightforward, but inadequate to the task: using a single characteristic velocity will result in a bouncingbarrier size at which dust grain collisions all result in bouncing. In reality, there is a range of collision velocities, and dust grains that collide at the low end can stick and grow (see Windmark 2012).

We analyzed the problem numerically, and generated analytical fits for inclusion in future dust coaggulation studies.

To enable our study of dust grain collisions in turbulence, we used artificial turbulence (velocity equation below), which side-stepped the resolution limitations.

$$V(x,t) = \sum_{m,n=0,1}^{m,n=8,3} \sqrt{2}a_{mn}(t)v_m \hat{v}_{mn} \cos \left[ k_{mn} \cdot x + \phi_{mn}(t) \right]$$

If interested, ask me what the terms above mean and why they were chosen.

#### **INERTIAL RANGE:**

**Dust grains** for which bouncing is expected to be important are too small/large to care about the largest/smallest scales of the turbulence: one needs a large inertial range, which is out of reach.

#### **IDENTICAL STOPPING TIMES (Hubbard 2012) DIFFERENT STOPPING TIMES (Hubbard 2013) Dust grains with identical stopping times are highly** correlated. Clustering We measure the stopping time ratio Velocity probability distribution with $\epsilon$ . Note that the mass ratio is Probability density **R**<sup>-0.52</sup> 100.00 Clustering High velocity collision (ε+1)<sup>3</sup>. 10.00 statistics depend only Three different separations 1.00 weakly on particle $\epsilon \equiv \frac{\tau_1}{-1} > 0$ 0.10 (log) scaled separation separation. 0.001 0.01 720.2 0.4 0.6 (Zoom) Low velocity collision Scaled collision speed statistics depend Zoom of low collision speeds extremely strongly **Dust grains** with *identical* stopping times 3.0



0.046

Clu dense enough clusters to trigger 1.5 the streaming instability of Johansen et. al. 2007. 0.5 1.0 2.0 2.5 1.5 3.0 Fitting collisional velocity probability distributions Different stopping times (see Hubbard 2013 for details)

σ

sterin

Clu

stering

2.0

0.000

3.0

2.5

2.0

**For ε>0.9** (about a 10-fold mass difference) the collisional velocity probability distribution is

By comparison, particles with different

**Clustering** is stronger for particles

**Unless the dust** size distribution



ε=0.0625

0.010

(linear) scaled separation

0.015

0.02



the collisional velocity probability

distribution for particle pairs with identical





**For ε<0.9** (about a 10-fold mass difference) the collisional velocity probability distribution is **not** Maxwellian.