



Wavelet-based cross-correlation analysis of the scaling in molecular clouds Tigran Arshakian & Volker Ossenkopf

Overview

We develop a wavelet-based cross-correlation (WCC) method to study the correlation between structural changes in molecular clouds as a function of scale. The method compares a pair of maps observed in different tracers or at different velocity ranges.

Advantages of the WCC method:

- Allows to measure the correlation coefficient and structural offset between two maps as a function of
- · Allows to weight individual pixels by their
- observational significance.

 Is robust against the noise.

Application of the WCC to simulated fBm (fractional Brownian motion) maps reveals that:

Cross-correlation coefficient can strongly depend on

- · Correlation coefficient and offset can be recovered robustly regardless of noise.

Analysis of the G333 molecular line maps (13CO and C¹⁸O) shows:

• A large scale gradient in the structural distribution.

This could indicate a density structure where every core shows a low density tail towards the South-West mainly seen in 13CO

The WCC can be used to trace the correlated structural changes between different maps of a molecular cloud at scales representing the structural and physical importance such as chemical and phase transitions

Wavelet cross-correlation (WCC)

Wavelet transform is proven to be a powerful tool for the scaling analysis in galaxies, interstellar clouds (Stutzki et al. 1998; Frick et al. 2001; Ossenkopf et al. 2008, hereafter O08). Convolution of the wavelet filter $\psi(l,r)$ (r = (x,y)) with an image f(r) filters the image on a scale l given by

$$F(\mathbf{r},l) = \int \int f(\mathbf{r})\psi(l,\mathbf{r})d\mathbf{r}$$

Optimal wavelet filter for molecular clouds is found to be the Mexican-hat filter (O08), which provides the correct power spectral slope and spectral

$$\psi(\mathbf{r}) = \psi_c(\mathbf{r}) + \psi_a(\mathbf{r}) = \frac{4}{\pi l^2} \exp\left(\frac{\mathbf{r}^2}{(l/2)^2}\right) + \frac{4}{\pi l^2(\nu^2 - 1)} \left[\exp\left(\frac{\mathbf{r}^2}{(\nu l/2)^2}\right) - \exp\left(\frac{\mathbf{r}^2}{(l/2)^2}\right) \right]$$

where $\psi_c(\mathbf{r})$ and $\psi_a(\mathbf{r})$ are the core and annulus of the filter, and v = 1.5 is

The Δ -variance (O08) of the filtered map F(r,l) and its weights $w_E(r,l)$, indicating the significance of each pixel of the image, shows characteristic scales and the power spectral slope in the individual maps.

To study the cross-correlation of the amount of structure between two maps, f and g, on different scales, we introduce the wavelet-based crosspower spectrum.

$$C_l(\tau) = \frac{C(l, \tau)}{\sigma_F(l) \, \sigma_G(l)} \tag{1}$$

$$C(l,\tau) = \iint_{x,y} \sqrt{w_F(\mathbf{r},l)w_G(\mathbf{r}+\tau,l)} F_w^*(\mathbf{r},l) G_w(\mathbf{r}+\tau,l) d\mathbf{r}$$

$$\sigma_F(l) = \left(\iint_{x,y} w_F(\mathbf{r}, l) \, F_w^2(\mathbf{r}, l) \, d\mathbf{r} \right)^{1/2}, \, \sigma_G(l) = \left(\iint_{x,y} w_G(\mathbf{r}, l) \, G_w^2(\mathbf{r}, l) \, d\mathbf{r} \right)^{1/2}$$

$$F_w = F(\mathbf{r}, l) - \overline{F}_w(l), \quad G_w = G(\mathbf{r}, l) - \overline{G}_w(l),$$

$$\overline{F}_w(l) = \frac{\iint w_F(\boldsymbol{r},l) F(\boldsymbol{r},l) \, d\boldsymbol{r}}{\iint w_F(\boldsymbol{r},l) \, d\boldsymbol{r}} \cdot \overline{G}_w(l) = \frac{\iint w_G(\boldsymbol{r},l) G(\boldsymbol{r},l) \, d\boldsymbol{r}}{\iint w_G(\boldsymbol{r},l) \, d\boldsymbol{r}}$$

are weighted means of two maps respectively

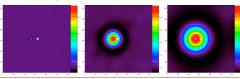
Cross-power spectrum (Eq. (1); see Fig 1) is used to estimate the correlation coefficient r(l) on scale l

$$r(l) = C_l(\tau = 0) \tag{2}$$

and **offset vector** $\tau(l) = [\Delta x, \Delta y]$, where the Δx and Δy give the location of the correlation peak of cross-power spectrum

$$r_{\max}([\Delta x, \Delta y], l) = \underset{x,y}{\operatorname{arg max}} C_l(\tau)$$

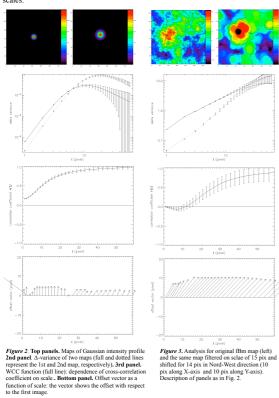
Figure 1. Cross-power spectrum of two fBm maps (see, for example, Fig. 3, top panel) for scales of 1 pix, 15 pix, and 30 pix. Offset between fBm maps is 10 pix along the X-axis. The distance between the peak of the maximum correlation and the map center (white cross) shows the offset $\tau(l)$ =10 pix between maps and the center of the cross-power spectrum gives the correlation coefficient $\tau(l)$ (see Eq. (2)).



Simple tests

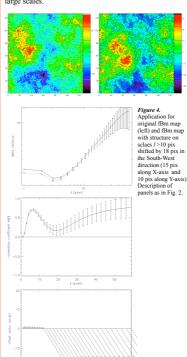
We start testing the WCC for two simulated circular structures having Gaussian intensity profiles with amplitude=1 and standard deviations of 3 pix and 5 pix, which are offset by 5 pix along Y-axis (Fig. 2, top panel). The size of the Gaussians is traced by the maxima in the Δ -variance spectra (10-15 pix; 2^{nd} panel). Correlation coefficient strongly depends on scale becoming large above the dominant structure scales. Amplitude and direction of the offset $\tau(l)$ (bottom panel) are correctly recovered on those

We also compared two fBm maps with spectral index of 3 (Fig. 3, top panel) where the second map is the filtered (maximum filter with a size of 15 pix) version of the first map, which mimics the opacity of optically thick lines. Correlation is negligible at small scales and it becomes significant at $l \ge 30$ pix (3rd panel). The correlation coefficient turns negative at scales below the mutual offset. The offset vector (length and direction) is exactly recovered for $l \ge 8$ pix (bottom panel)



Realistic tests

In reality the offset can vary on different spatial scales. To test this, we generate a fBm map with S/N=10 (Fig. 4 ton left panel) and use the Fourier shift theorem to shift the large scales ($l \ge 10$ pix) by 18 pix in the South-West direction (top right panel). The Δ -variance spectrum has a bump at small scales due to low S/N value (2nd panel). Application of the WCC for these two maps shows that the offset vector can be recovered robustly (bottom panel). Structures are correlated significantly at small scales ($l \le 10$ pix) as expected from the construction of the maps. The correlation gradually decreases from 10 to 20 pix (3rd panel) due to the shift of a structure at scales l > 10 pix. At scales of \geq 20 pixels the structure sizes exceed the offset so that we find a gradual increase of the correlation coefficient at large scales



Application to GMC G333

We applied the WCC to the maps of the giant molecular cloud (GMC) G333 observed in ¹³CO and C¹⁸O 1-0 (Lo et al. 2009; Fig. 5, top panels). Their Δ -variance spectra have pronounced structure on scales of about 30 pix and 50 pix respectively and the structure is strongly correlated at scales > 5 pix (3rd panel). Large differences between the maps only occur at small scales affected by noise. The structure is offset at scales larger than 40 pix with amplitude of ~ 5 pix along the filament, where all structures in 13CO are shifted to the South-West relative to C18O. This could indicate a large-scale column density gradient enhancing the structure in 13CO in that direction

